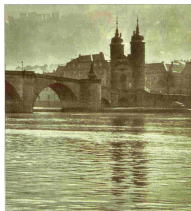


Gauge Data from Chow Groups and Massless Matter in F-Theory

joint work with M. Bies, C. Pehle and T. Weigand:
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- ▶ To obtain full massless matter spectrum, need gauge data beyond four-form flux;
- ▶ Chow groups will give us handle on them;

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- ▶ Supersymmetry: $G_4 \in H^{2,2}(Y_4)$;
- ▶ Has to be quantised: $G_4 + \frac{c_2}{2} \in H^4(Y_4, \mathbb{Z})$;

Chirality

- ▶ Type IIB: chirality along curve of intersecting branes given by

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- ▶ Matter surfaces, C_{R_q} , consist of linear combinations of blow-up \mathbb{P}^1 's fibred over enhancement curve C_{R_q} ;
- ▶ Recall: linear combination such that in dual M-theory picture, M2-brane wrapping this combination is one state of R_q ;

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- ▶ Splitting of Pic encoded via:

$$0 \rightarrow \mathcal{J}^1(X) \rightarrow \text{Pic}(X) \xrightarrow{c_1} H_{\mathbb{Z}}^{1,1}(X) \rightarrow 0$$

$c_1(L) = \frac{1}{2\pi}F$ linear map onto $H^{1,1}(X, \mathbb{Z})$,
 $\mathcal{J}^1(X) = H^{0,1}(X, \mathbb{C})/H^1(X, \mathbb{Z})$ Jacobian of X (space of flat connections).

Beyond the Chiral Index II

- ▶ For A_3 & G_4 in F-theory exists similar decomposition:

$$0 \rightarrow \underbrace{\mathcal{J}^2(\hat{Y}_4)}_{\substack{2^{\text{nd}} \text{ intermediate} \\ \text{Jacobian}}} \rightarrow \underbrace{H_D^4(\hat{Y}_4, \mathbb{Z}(2))}_{\substack{4^{\text{th}} \text{ Deligne} \\ \text{cohomology class}}} \xrightarrow{c_2} H_{\mathbb{Z}}^{2,2}(\hat{Y}_4) \rightarrow 0$$

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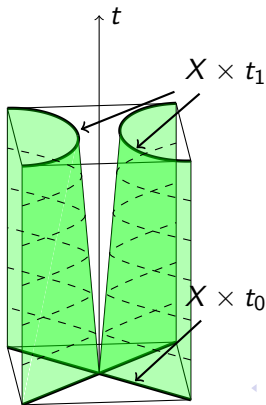
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- ▶ But can work indirectly by using Chow groups;

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- ▶ Rational equivalence: $C_1 \cong C_2 \in Z_n(X)$ if $C_1 - C_2$ is zero/pole of meromorphic function defined on $(n+1)$ -dim. irreducible subvariety of X ; Equivalently: two algebraic cycles $C_1, C_2 \in Z_i(X)$ rationally equivalent if \exists rationally parametrised family of cycles interpolating between them;



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 & & \downarrow AJ & & \downarrow \hat{\gamma}_p & & \downarrow & & \\
 0 & \longrightarrow & J^p(X) & \longrightarrow & H_D^{2p}(X, \mathbb{Z}(p)) & \xrightarrow{\hat{c}_p} & H_{\mathbb{Z}}^{p,p}(X) & \longrightarrow & 0
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 2. Manipulations modulo rational equivalence preserve C_3 modulo gauge equivalence;

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- ▶ This collection of points $A_{R,G} \in Z_0(\mathcal{C}_R)$ can be used to define line bundle $L_{G,R} = \mathcal{O}_{\mathcal{C}_R}(A_{R,G})$ on \mathcal{C}_R ;

Matter II

► **Proposal:**

[Bies,CM,Pehle,Weigand]

massless $\mathcal{N} = 1$ chiral multiplets counted by

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- ▶ Can checked proposal for fluxes/gauge data coming from e.g. $U(1)$ -symmetries;

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⇒ $\gamma(w_A)$ together with $f \in \text{CH}^1(B_3)$ gives four-form flux (class)

$$G_4^A = \pi^* \gamma(f) \cup \gamma(w_A) \in H^{2,2}(\hat{Y}_4)$$

which satisfies 'one leg ...' condition and leaves all non-abelian sym. untouched.

Applied to U(1)-model II

- ▶ Have now more than flux, because

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- ▶ In many cases, CY four-fold embedded in toric ambient space and w_A is pullback ($w_A = j^* \tilde{w}_A$) then

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Applied to $U(1)$ -model III

- ▶ In such cases, can use intersections of toric ambient variety to calculate $\alpha_R \cdot j_{\iota_R} \tilde{w}_A$ and find:

$$\pi|_{C_R*}(\alpha_R \cdot \iota_R w_A) = \pi|_{C_R*}(\alpha_R \cdot j_{\iota_R} \tilde{w}_A) = q_A(R)[C_R], \in \text{CH}_1(C_R)$$

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- ▶ Finally from $[C_R] \cdot \iota_{R|B_3} f$ obtain collection of points $A_{R,A} \in Z_0(C_R)$ on C_R ;

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⇒ Defines line bundle $\mathcal{O}_{C_R}(A_{R,A})$; Massless matter states correspond to

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- ▶ By means of *cohomCalc* and self-written Mathematica code, could work out spectral sequences to obtain cohomologies of line bundles on curves;

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curve	$h^0(C, \mathcal{L} _C)$	representation	$h^1(C, \mathcal{L} _C)$	representation
C_{10}	4	$\mathbf{10}_{-1}$	1	$\overline{\mathbf{10}}_{+1}$
$C_{\bar{5}_m}$	6	$\overline{\mathbf{5}}_3$	3	$\mathbf{5}_{-3}$
C_{5_H}	9	$\mathbf{5}_2$	9	$\overline{\mathbf{5}}_{-2}$
C_1	585	$\mathbf{1}_5$	0	$\overline{\mathbf{1}}_{-5}$

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- ▶ Apply our methods to cases where intersections have to be done on CY itself;

Thank you for your attention!